Measurement of gamma-rays from the keV-neutron capture by $^{139}$La

M. Igashiraa, M. Saito, J. Nishiyama, T. Ohsakib, and T. Katabuchi

Research Laboratory for Nuclear Reactors, Tokyo Institute of Technology, Tokyo 152-8550, Japan

Abstract. The capture cross sections and capture γ-ray spectra of $^{139}$La were measured in a neutron energy region from 10 to 100 keV and at 550 keV. Pulsed keV neutrons were produced by the $^7$Li(p,n)$^7$Be reaction with a pulsed proton beam from a Pelletron accelerator. Capture γ-rays were detected with a large anti-Compton NaI(Tl) spectrometer by means of a time-of-flight method. A pulse-height weighting technique was applied to the observed capture γ-ray pulse-height spectra to obtain capture yields. Using the standard capture cross sections of $^{197}$Au, the capture cross sections of $^{139}$La were derived with uncertainties of about 5%. The capture γ-ray spectra were obtained by unfolding the observed capture γ-ray pulse-height spectra. The present results were compared with previous measurements and evaluated values.

1 Introduction

Since the neutron magic nuclides have small keV-neutron capture cross sections, nuclides with $N = 50, 82,$ and $126$ act as bottlenecks in the chain of neutron capture reactions in the s-process nucleosynthesis. Therefore, their accurate keV-neutron capture cross sections are important for the study on the s-process nucleosynthesis.

On the other hand, neutron capture γ-ray spectra include important information on physical quantities such as γ-ray strength function. Therefore, simultaneous analysis of both cross sections and γ-ray spectra will provide important information on the keV-neutron capture reaction mechanism.

From these viewpoints, we have measured the keV-neutron capture cross sections and capture γ-ray spectra of $^{139}$La ($N = 82$). In the present paper, we report on the experimental results.

2 Experimental procedure

The experimental procedure has been described in detail elsewhere [1], so it is summarized here.

The capture cross sections and capture γ-ray spectra of $^{139}$La were measured in an incident neutron energy region from 10 to 100 keV and at 550 keV, using the 3-MV Pelletron accelerator of the Research Laboratory for Nuclear Reactors at the Tokyo Institute of Technology. An experimental arrangement is shown in figure 1.

Pulsed neutrons were produced by the $^7$Li(p,n)$^7$Be reaction with a pulsed proton beam (1.5 ns width, 4 MHz repetition rate) from the Pelletron accelerator. The incident neutron energy spectrum on a capture sample was measured by means of a time-of-flight (TOF) method with a $^6$Li-glass scintillation detector. A 5.0-mm diameter by 5.0-mm $^6$Li-glass detector located 30 cm from the neutron source was used for the measurement in the 10–100 keV region, and a 102.0-mm diameter by 6.4-mm $^6$Li-glass detector located 460 cm was used for that at 550 keV.

The capture sample was natural La$_2$O$_3$ powder (about 50 g) which was contained in a graphite case with an inner diameter of 55 mm and an inner length of 14 mm. A gold (Au) disk with a diameter of 55 mm and a thickness of 1 mm was used as a standard sample. Each capture sample was located at an angle of 0 degree with respect to the proton beam direction. The distance between the neutron source and the sample was 12 cm for the measurement in the 10–100 keV region and 20 cm for that at 550 keV.

The capture γ-rays emitted from the $^{139}$La or $^{197}$Au sample were measured with a large anti-Compton NaI(Tl) spectrometer by means of a TOF method. The main NaI(Tl) detector of the spectrometer had a diameter of 15.2 cm and a length of 30.5 cm, and was centered in a hollow Compton-suppression NaI(Tl) detector with an outer diameter of 33.0 cm and a length of 35.6 cm. The spectrometer was set in a heavy shield consisting of borated paraffin, borated polyethylene, Cd, $^6$LiH and potassium free lead [2]. The capture γ-rays were observed at an angle of 125 degree with respect to the proton beam direction. The signals from the spectrometer were recorded in a personal computer as two-dimensional data of pulse-height (PH) and TOF.

The measurements with the $^{139}$La sample (La run), the $^{197}$Au sample (Au run), and...
no sample (Blank run) were performed cyclically to average out changes in experimental conditions such as the incident neutron spectrum.

3 Data processing

The data processing method has been also given in detail elsewhere [1], so it is summarized here.

The incident neutron energy spectrum on the sample in the 10–100 keV or 550 keV measurement was derived from the TOF spectrum observed with the small or large $^6$Li-glass detector in Blank run, respectively. The normalized spectra are shown in figure 2.

Figure 3 shows the TOF spectra observed with the $\gamma$-ray spectrometer in the 10–100 keV measurement. The broad peak below 550 channel in the spectrum for La run or Au run is due to the capture $\gamma$-rays from the corresponding sample, and the sharp peak around 650 channel is due to the $\gamma$-rays from the $^7$Li($p,\gamma)^8$Be reaction in the Li target. Four gates were set in the foreground region and one gate was set in the time-independent background region, as shown in figure 3, and the corresponding PH spectra were obtained from the two-dimensional data. Then, the net capture $\gamma$-ray PH spectra were obtained by subtracting the background PH spectra normalized with the gate-width ratios from the foreground PH spectra.

In order to obtain the capture yields corresponding to the individual gates of La and Au runs, a PH weighting technique [3] was applied to the net capture $\gamma$-ray PH spectra. The number of incident neutrons in Au run was determined by the capture yield of $^{197}$Au and the averaged capture cross section of $^{197}$Au, which was obtained from the capture cross sections of ENDF/B-VI [4] and the normalized neutron energy spectrum. The number of incident neutrons in La run was derived from that in Au run and the neutron monitor counts of the $^6$Li-glass detector. The averaged neutron capture cross section of $^{139}$La was derived from the number of incident neutrons and the capture yield of the $^{139}$La sample.

Corrections were made for the neutron self-shielding and multiple-scattering in the sample [5], for the $\gamma$-ray scattering and absorption in the sample, for the effect of the discrimination level (600 keV) in the net capture $\gamma$-ray PH spectra on the capture yields, and for the effect of impurities in the sample on the capture yields.

The capture $\gamma$-ray spectra were derived by unfolding the net capture $\gamma$-ray PH spectra with a computer code, FERDOR [6], and the response matrix of the $\gamma$-ray spectrometer. The derived capture $\gamma$-ray spectra were normalized so that the integration of the spectrum with respect to the $\gamma$-ray energy is equal to the neutron binding energy plus the incident neutron energy [1].

4 Results and discussion

The capture cross sections of $^{139}$La were derived with uncertainties of about 5% in the incident neutron energy region from 10 to 100 keV and at 550 keV. The results in the 10–100 keV region are shown in figure 4, and compared with
previous measurements and the evaluations of JENDL-3.3 [7] and ENDF/B-VI [8]. Since activation methods are applicable to the measurement of capture cross sections of $^{139}$La, there are a considerable number of previous experimental results by activation methods. However, those results are not shown in figure 4.

The results of Musgrove et al. [9] are in good agreement with the present results. On the other hand, the result at 30 keV of Macklin et al. [10] is larger than the present one by 25%, and their result at 65 keV is smaller than the present one by about 25%. As for the results of Konks et al. [11], the neutron energy dependency of their results is different from that of the present results, although their result around 40 keV is consistent with the present results.

The differences between both evaluations are less than 15% in the neutron energy region from 20 to 100 keV, and the presents results lie between both evaluations. It is worthy of note that the evaluations of ENDF/B-VII [12] in the region from 20 to 100 keV are almost the same as those of JENDL-3.3.

The derived capture $\gamma$-ray spectrum of $^{139}$La at 550 keV is shown in figure 5, and compared with the result of Voignier et al. [13] at 500 keV. Low lying states of the residual nucleus, $^{140}$La, are shown in figure 5. In the present spectrum, strong transitions from the neutron capture state to low lying states are clearly observed, although those transitions are not observed in the spectrum of Voignier et al. It is worthy of note that the data below 1.5 MeV in the spectrum of Voignier et al. are not observed ones but calculated ones with a statistical model.

5 Conclusion

The neutron capture cross sections and capture $\gamma$-ray spectra of $^{139}$La were measured in an energy region from 10 to 100 keV and at 550 keV. The capture cross sections were derived with uncertainties of about 5%. The derived cross sections in the region from 10 to 100 keV lay between the JENDL-3.3 (and ENDF/B-VII) and ENDF/B-VI evaluations. The derived capture $\gamma$-ray spectra at 550 keV was quite different from that of Voignier et al. at 500 keV for the transitions from the neutron capture state to low-lying states of the residual nucleus.

References

7. JENDL-3.3 data file for $^{139}$La (MAT = 5728), evaluated by T. Aoki et al. (1993).
8. ENDF/B-VI data file for $^{139}$La (MAT = 5728), evaluated by R.E. Schenter, F. Schmittroth (1980).