

Neutron capture measurements on Tl-isotopes at DANCE

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Abstract. The thallium isotopes play an important role in the *s*-process nucleosynthesis at the *s*-process endpoint. Furthermore, ²⁰⁴Tl is one of few branch point isotopes in the endpoint region. The understanding of branch point isotopes provides modeling constraints on the temperatures and neutron densities during which the process takes place. The production of *s*-only ²⁰⁴Pb is controlled almost entirely by ²⁰⁴Tl. Measurements of the capture cross-sections of the stable Tl isotopes have recently been made using the DANCE 4 π array at LANSCE. This provides needed resonance information in the region as well as preparing the way for measurements of as yet unmeasured capture cross-section of the unstable ²⁰⁴Tl.

1 Introduction

About 50% of the element abundances beyond iron are produced via slow neutron capture nucleosynthesis (*s* process). Starting at iron-peak seed, the *s*-process mass flow follows the neutron-rich side of the valley of stability. If different reaction rates for the same isotope are comparable, the *s*-process path branches and the branching ratio reflects the physical conditions in the interior of the star. Such nuclei are most interesting, because they provide the tools to effectively constrain modern models of the stars where the nucleosynthesis occurs. As soon as the β^- decay is faster than the typically competing neutron capture, no branching will take place. Thus experimental neutron capture data for the *s* process are only needed if the respective neutron capture time under stellar conditions is similar to or shorter than the β^- decay time; this includes all stable isotopes. Depending on the actual neutron density during the *s* process, the "line of interest" is closer to or farther away from the valley of stability. Figure 1 shows a summary of the neutron capture and β^- decay times for radioactive isotopes on the neutron rich side of the valley of stability, under the condition that the classical neutron capture occurs faster than the terrestrial β^- decay. The vast majority of isotopes where an experimental neutron capture cross section is desirable have β^- half-lives of at least hundreds of days. Such isotopes can be investigated with the DANCE array.

Since the neutron densities during the *s* process are not high enough to overcome the α -unstable isotopes beyond ²⁰⁹Bi, the Tl-Pb-Bi group marks the termination region of the *s*-process (fig. 2). The production of the *s*-only ²⁰⁴Pb is mostly determined by the branching at ²⁰⁴Tl. Another interesting branching occurs at ²⁰⁵Pb, especially since ²⁰⁵Tl, which is stable under laboratory conditions, becomes unstable under stellar conditions [2,3]. The stellar site for the production

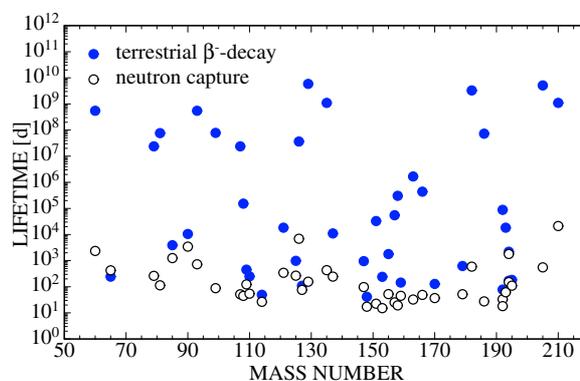


Fig. 1. Neutron capture life times (open circles) and terrestrial β^- life times (filled circles) for unstable isotopes on the classical *s*-process path as a function of mass number. Shown are only isotopes where the neutron capture is faster than the stellar β^- decay for a neutron density of $5 \cdot 10^8 \text{ cm}^{-3}$ at a temperature of 30 keV. The neutron capture cross sections are mostly taken from [1] and the stellar decay rates from [2].

of these isotopes are metal-poor Asymptotic Giant Branch (AGB) stars [4,5]. Both radioactive isotopes are very difficult to produce under terrestrial conditions. Irradiating Bi samples in close proximity of a spallation source and chemically separating the freshly produced material from the bismuth would be a very good solution, especially for ²⁰⁴Tl. Further enrichment by mass separation is not necessary, since ²⁰⁴Tl has the highest neutron capture Q-value of all relevant Tl isotopes. The neutron capture measurements on ^{203,205}Tl described in this article offer a significant improvement in the neutron capture data on the stable isotopes of thallium, which is important since the present data are limited in resolution and the resolved resonance region reaches above 10 keV, particularly for ²⁰⁵Tl. In addition, the present measurements are in preparation for a ²⁰⁴Tl(*n*, γ) measurement at a later stage.

2 The DANCE array

The Detector for Advanced Neutron Capture Experiments (DANCE) is designed as a high efficiency, highly segmented

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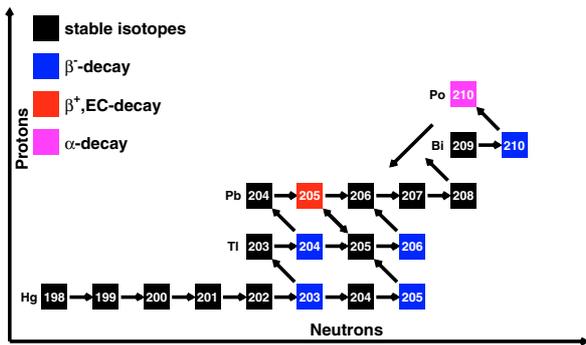


Fig. 2. The s-process path at its termination region between mercury and polonium.

4π BaF₂ detector for calorimetrically detecting γ -rays following a neutron capture. DANCE is located on the 20 m neutron flight path 14 (FP14) at the Manuel Lujan Jr. Neutron Scattering Center at the Los Alamos Neutron Science Center (LANSCE) [6]. The initial design work is described in [7]. For practical reasons the detector modules do not cover the entire solid angle. The design of the detector is such that a full 4π array would consist of 162 crystals of four different shapes, each shape covering the same solid angle [8]. Two of the 162 crystals are left out in order to leave space for the neutron beam pipe. Depending on the experiment, one crystal can be replaced by a sample changer mechanism, which makes it possible to exchange up to 3 samples without closing the beam shutter and breaking the vacuum of the beam pipe. Thus the full array is designed to host 159 or 160 out of 162 possible BaF₂ crystals. The dimensions of the bare crystals are designed to form a BaF₂ shell with an inner radius of 17 cm and a thickness of 15 cm. Thanks to the fairly low repetition rate of 20 Hz, measurements can be carried out over the whole energy range from 10 meV to 500 keV. This combination of a strong neutron source and a high efficiency γ -ray detector allows to measure (n, γ) cross section of radioactive isotopes down to a few hundred days half-life. Further details on the overall performance of the array can be found in [9, 10].

3 Samples

Altogether we used four thallium samples. Since the cross section changes over several orders of magnitude over the investigated energy range, we used two ²⁰³Tl samples with 7.5 mg and 51 mg and two ²⁰⁵Tl samples with 9.5 mg and 53 mg. The smaller samples were dried out of a solution onto a thin mylar foil, while the larger ones were small pieces of thallium mounted on a thin mylar foil. Therefore the amount of material per area was much lower for the small samples, hence they were much better suited for analyzing the strong resonances. All of the samples were sealed with VYNS to prevent the thallium from migrating or flaking. Table 1 contains the details for each of the samples.

4 Experiment

The data analysis is still in progress. Therefore only preliminary data will be presented here. In total, about 30 days of beam time were spent on ^{203,205}Tl, ¹⁹⁷Au, and background runs.

Table 1. Masses and isotopic composition of the four samples used during the experiment.

sample nbr.	mass (mg)	²⁰³ Tl (%)	²⁰⁵ Tl (%)
1	51	97.15 ± 0.04	2.85
2	7.5	97.15 ± 0.04	2.85
3	53	0.26	99.74 ± 0.01
4	9.5	0.26	99.74 ± 0.01

5 Preliminary results

5.1 ²⁰³Tl

The experimental information on the ²⁰³Tl(n, γ) reaction above 1 keV is sparse (see fig. 3). Based on the measurements and evaluation by Macklin and Winters [11], the recommended value for the Maxwellian averaged cross section (MACS) at 30 keV by Bao et al. is 124 ± 8 mb [1].

Figure 4 shows the response of the DANCE array during the experiment for low neutron energies. Shown is the total

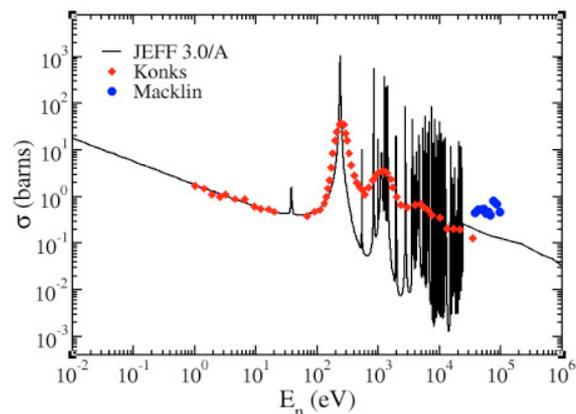


Fig. 3. Experimental [11, 12] data as reported at EXFOR and evaluated [13] data for the ²⁰³Tl(n, γ) reaction in the higher energy region before this measurement. Additionally resonance parameters between 3 and 13 keV are reported [11].

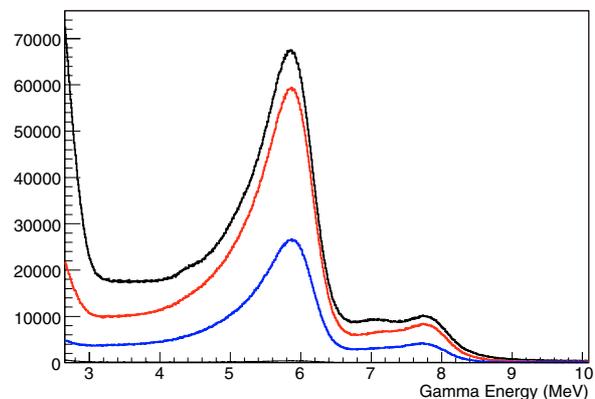


Fig. 4. Total energy deposited in the DANCE array for the ²⁰³Tl sample for low neutron energies between 10 meV and 1 keV. The different curves correspond to different multiplicities: Multiplicity 1–3 from top to bottom.

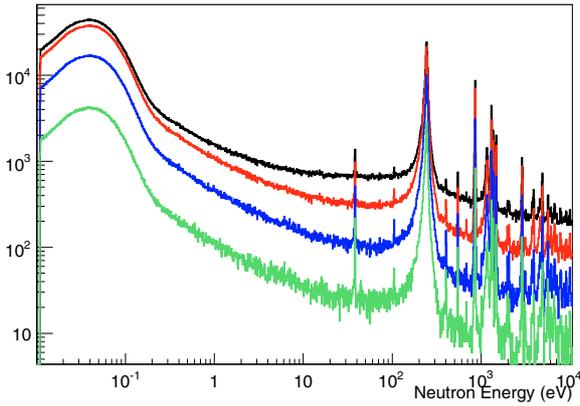


Fig. 5. Time-of-flight spectrum for ^{203}Tl with an cut on the total deposited energy between 5 and 6.5 meV. The different curves correspond to multiplicities 1 (top) to 4 (bottom).

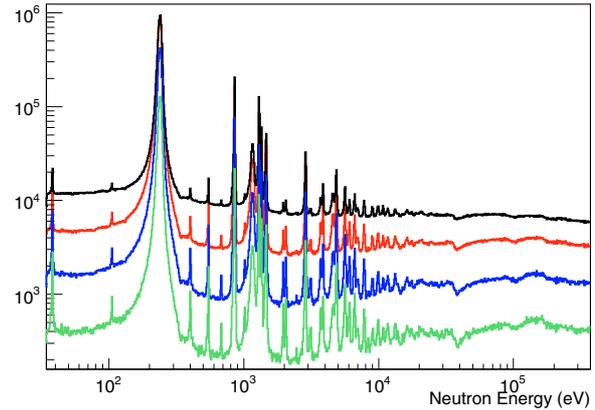


Fig. 7. Time-of-flight spectrum for ^{203}Tl and multiplicities 1 to 4 (top to bottom).

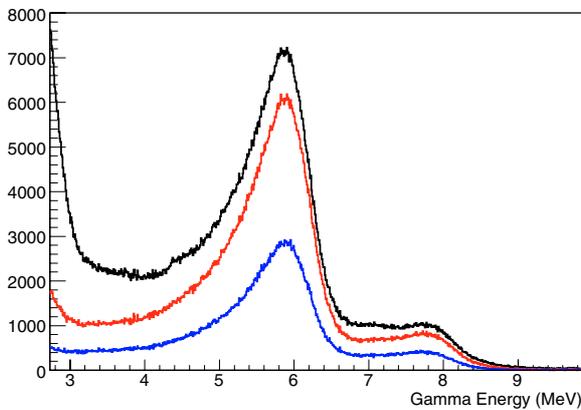


Fig. 6. Total energy deposited in the DANCE array for the ^{203}Tl sample for high neutron energies between 30 eV and 200 keV. The different curves correspond to different multiplicities: Multiplicity 1-3 from top to bottom.

energy deposited in the detector as a function of cluster multiplicity, which is very close to the multiplicity of the emitted γ -rays [7,9]. The Q-value of the reaction is 6.65 meV, resulting in the pronounced peaks between 5 and 6.5 meV. A gate on these peaks improves the signal-to-background ratio. The events above 7 meV result from captures of scattered neutrons in the BaF_2 crystals.

Figure 5 shows results of a total-energy cut on time-of-flight (TOF) spectrum of different multiplicities. The signal-to-background improves with increasing multiplicity since most background components tend to yield low multiplicities [9]. The trade-off typically is reduced statistics.

Figure 6 shows a similar spectrum as figure 4, but for higher neutron energies. The capture-to-scatter ratio is typically decreasing with increasing neutron energy, which makes neutron capture measurements increasingly difficult. This effect is not very pronounced for $^{203}\text{Tl}(n,\gamma)$, but can still be observed.

Figure 7 shows a TOF spectrum for the events shown in figure 6 with an additional total-energy cut between 5 and 6.5 meV. Data have been successfully taken up to 200 keV neutron energy.

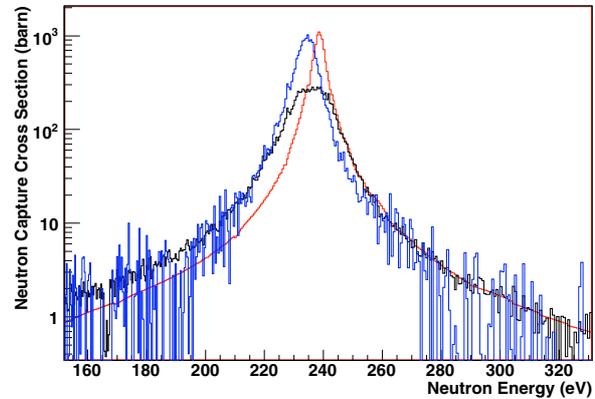


Fig. 8. Neutron capture cross section for ^{203}Tl derived from DANCE (black, blue) (multiplicity cut from 1 to 4 and total-energy cut from 5 to 6.5 meV) and a comparison with evaluated data (red) [13] for the biggest resonance in the $^{203}\text{Tl}(n,\gamma)$ cross section (see fig. 6). The data for the black curve were taken with sample 1 (see table 1) and are not corrected for self-absorption, while the data for the blue curve were taken with the much thinner sample 2.

Figure 8 shows the conversion from TOF spectrum to cross section in the neutron energy region around the biggest resonance of the $^{203}\text{Tl}(n,\gamma)$ reaction. The plot illustrates the advantage of using samples of different thickness. Significant self-absorption corrections would be necessary for sample 1, while sample 2 reproduces the shape of the resonance as suggested by the JEFF-3.0/A evaluation [13]. The final data will be derived from the smaller sample 2 for the biggest resonances and from the bigger sample 1 elsewhere (see also table 1).

5.2 ^{205}Tl

Similar to the situation for ^{203}Tl , the experimental information on the $^{205}\text{Tl}(n,\gamma)$ reaction above 1 keV is sparse (see fig. 9). Based on the measurements and evaluation by Macklin and Winters [11], the recommended value for the Maxwellian averaged cross section (MACS) at 30 keV by Bao et al. is 54 ± 4 mb [1].

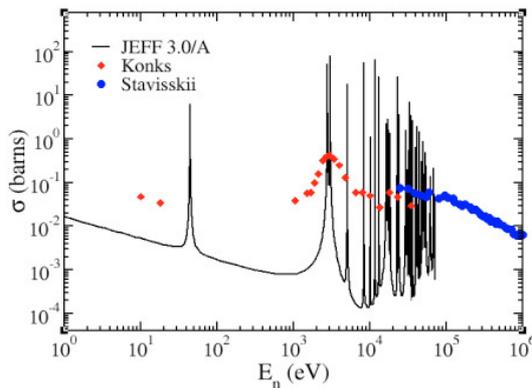


Fig. 9. Experimental [12, 14] data as reported at EXFOR and evaluated [13] data for the $^{205}\text{Tl}(n,\gamma)$ reaction before this measurement.

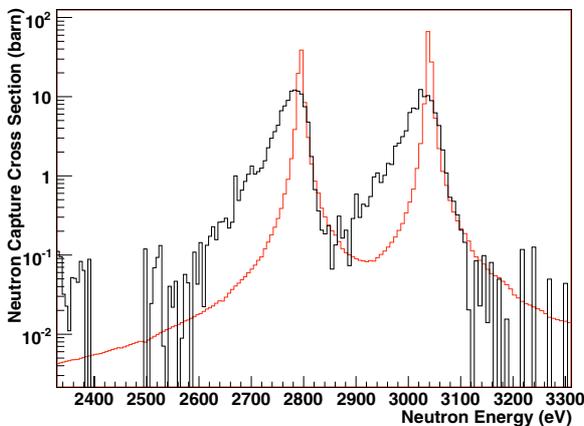


Fig. 10. Neutron capture cross section for ^{205}Tl derived from DANCE (black) (multiplicity cut from 2 to 3 and total-energy cut from 5 to 6.5 meV) and a comparison with evaluated data (red) [13] for the biggest resonance doublet in the $^{205}\text{Tl}(n,\gamma)$ cross section (see fig. 9).

Figure 10 shows the conversion of the data to cross section in the neutron energy region around the resonance doublet of the $^{205}\text{Tl}(n,\gamma)$ reaction. The plot illustrates the improved accuracy of the DANCE data compared to the data reported so far (see fig. 9). As for $^{203}\text{Tl}(n,\gamma)$, the final data will be derived from the smaller sample 4 for the biggest resonances and from the bigger sample 3 elsewhere.

6 Conclusion

Neutron capture data on enriched samples of the stable isotopes of thallium have been taken with the DANCE array at the Los Alamos National Laboratory. The goal of the experiment

was twofold. Firstly, we wanted to improve the available experimental data in the astrophysically interesting energy region between 100 eV and 200 keV. Secondly, we wanted to prepare for the neutron capture measurement on the important s -process branch point ^{204}Tl .

The preliminary data presented in this article suggest that both goals can be achieved as soon as the final data analysis has been performed.

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References

1. Z.Y. Bao, H. Beer, F. Käppeler, F. Voss, K. Wisshak, T. Rauscher, *At. Data Nucl. Data Tables* **76**, 70 (2000).
2. K. Takahashi, K. Yokoi, *At. Data Nucl. Data Tables* **36**, 375 (1987).
3. K. Yokoi, K. Takahashi, M. Arnould, *A&A* **145**, 339 (1985).
4. U. Ratzel, C. Arlandini, F. Käppeler, A. Couture, M. Wiescher, R. Reifarh, R. Gallino, A. Mengoni, C. Travaglio, *Phys. Rev. C* **70**, 065803 (2004).
5. C. Arlandini, F. Käppeler, K. Wisshak, R. Gallino, M. Lugaro, M. Busso, O. Straniero, *ApJ* **525**, 886 (1999).
6. P.W. Lisowski, C.D. Bowman, G.J. Russell, S.A. Wender, *Nucl. Sci. Eng.* **106**, 208 (1990).
7. M. Heil, R. Reifarh, M.M. Fowler, R.C. Haight, F. Käppeler, R.S. Rundberg, E.H. Seabury, J.L. Ullmann, J.B. Wilhelm, K. Wisshak et al., *Nucl. Instrum. Meth. A* **459**, 229 (2001).
8. D. Habs, F.S. Stephens, R.M. Diamond, *Tech. Rep., Report LBL-8945*, Lawrence Berkeley Laboratory (1979).
9. R. Reifarh, T.A. Bredeweg, A. Alpizar-Vicente, J.C. Browne, E.I. Esch, U. Greife, R.C. Haight, R. Hatarik, A. Kronenberg, J.M. O'Donnell et al., *Nucl. Instrum. Meth. A* **531**, 528 (2004).
10. J.L. Ullmann, U. Agvaanlvsan, A. Alpizar-Vicente, E. Bond, T.A. Bredeweg, E.I. Esch, C.M. Folden, U. Greife, R. Hatarik, R.C. Haight et al., *AIP Conf. Proc.* **769**, 918 (2005).
11. R.L. Macklin, R.R. Winters, *ApJ* **208**, 812 (1976).
12. V.A. Konks, F.L. Shapiro, *Soviet Physics JETP-USSR* **20**, 531 (1965).
13. NEA/EFF working group, *Tech. Report JEFF Report 19*, OECD Nuclear Energy Agency, www.nea.fr/html/dbdata/nds_jefreports/jefreport-19/jefreport-19.pdf (2005).
14. Y.Y. Stavisskii, V.A. Tolstikov, *Reactor Science and Technology (J. Nucl. Energy Parts A and B)* **16**, 496 (1962).